

Search for extraterrestrial neutrino-induced cascades using IceCube 79-strings

THE ICECUBE COLLABORATION¹ ¹See special section in these proceedings

Abstract: IceCube, a cubic kilometer detector at the South Pole, is the largest neutrino telescope currently taking data. Utilizing the transparent ice of Antarctica as a detection medium, IceCube digital optical sensors observe Cherenkov radiation from secondary particles produced in neutrino interactions inside or near the detector. Charged current v_{μ} interactions create muon tracks, while charged current v_e interactions, and neutral current interactions of all flavors initiate electromagnetic and hadronic showers (cascades). The background coming from atmospheric muons and muon bundles is many orders of magnitude larger than the cascade signal and makes it difficult to observe cascades. However, cascades have better energy resolution and lower atmospheric background compared to track-like events. The energy spectrum of extraterrestrial neutrinos is expected to be harder than that of atmospheric neutrinos. Thus using cascade events to search for a hardening of the energy spectrum is advantageous compared to using muon tracks. The search for extraterrestrial neutrino-induced cascades with energies in the tens of TeV to a few PeV neutrino energy range using improved reconstruction methods will be presented. The analysis uses 317 days of livetime of the data taken from May 2010 to May 2011 when 79 IceCube strings were operational. The analysis method and new results from the fully unblinded dataset will be presented.

Corresponding authors: Mariola Lesiak-Bzdak² (*mlesiak-bzdak@icecube.wisc.edu*), Achim Stöeßl³ (*achim.stoessl@desy.de*)

² Department of Physics and Astronomy, Stony Brook University. Stony Brook, NY 11794-3800

33

³ DESY, Platanenallee 6, 15738 Zeuthen, Germany

Keywords: IceCube, Neutrino, Cascades, Diffuse

1 **1 Introduction**

Extraterrestrial neutrinos, anticipated to be produced to-34 2 gether with cosmic rays, might provide information about 35 3 the mechanism of cosmic ray production and help to unveil 36 4 cosmic ray sources. Although neutrino fluxes from such 37 5 sources could be too low to be measured individually, an $_{38}$ 6 integrated flux over all sources might be possible to detect 7 with IceCube [1], a cubic kilometer scale neutrino telescope 39 8 located at the geographic South Pole. Incoming neutrinos 40 9 interact mostly via deep-inelastic nucleon scattering and 41 10 produce showers of secondary charged particles. Having 42 11 relativistic velocities, these particles produce Cherenkov₄₃ 12 light that is detected by Digital Optical Modules (DOMs). 44 13 Neutrino-nucleon reactions are induced by all neutrino 14 flavors via neutral current (NC) or charged current (CC) 45 15 interactions. In charged current reactions the charged lepton ⁴⁶ 16 is produced, which carries on average 50% (for $E_v \sim 10^{47}$ 17 GeV) to 80% (at high energies) of the neutrino energy;48 18 the remainder of the energy is transferred to the nuclear $_{49}$ 19 target. Depending on the charged lepton created in CC_{50} 20 reactions, neutrino flavor specific hit-patterns might be ⁵⁰ observed in the detector which allow the identification of the ⁵¹ 21 22 incoming neutrino flavor. Charged current ν_{μ} interactions ⁵² 23 create track-like hit patterns while CC v_e reactions produce ⁵³ 24 an electromagnetic and hadronic cascade which yields a 54 25 spherical hit-pattern. The typical cascade analysis searches 55 26 for v_e and v_{τ} from CC and all neutrino flavors from NC ₅₆ 27 interactions. 28

A previous cascade analysis searching for an astrophysi-⁵⁷ cal neutrino flux in IceCube with 22-strings instrumented ⁵⁸ [2] set a limit of 3.6×10^{-7} GeV· sr⁻¹s⁻¹cm⁻² at 90% C.L.⁵⁹ on E^{-2} astrophysical neutrinos (assuming a 1:1:1 flavor ra-60 tio) with 90% of events in the energy range between 24 TeV to 6.6 PeV. Another IceCube cascade analyses looking for an extraterrestrial neutrino signal using 40 strings obtained preliminary results [3] and set a limit at 90% confidence level on an astrophysical neutrino flux of 9.5×10^{-8} GeV· sr⁻¹s⁻¹cm⁻² with 90% of events in the energy range between 89 TeV to 21 PeV [4]. Preliminary cascade results using the 59-string configuration of IceCube were recently obtained and are presented at this conference [5].

In the recent 79- and 86-string IceCube detector, searches for extremely-high energy (EHE) neutrinos from all flavors from CC and NC interactions, two neutrino-induced cascade events at energies of 1 PeV were observed [6]. As a followup analysis, an all-sky search for all flavor neutrino events from CC and NC interactions with energies $E_v > 100$ TeV and neutrino first interaction well contained in the 79and 86-string IceCube detector, was performed and the preliminary results are presented in these proceedings [7].

The analysis described here was developed using Monte Carlo simulation and searched for an E^{-2} astrophysical neutrino-induced cascade flux within IceCube with 79 strings instrumented. In these proceedings, we present all flavor sensitivity using high-energy contained cascade events in the IceCube detector. We also discuss adding partially contained events, to increase the effective volume. The neutrino energy range in this analysis is between 44 TeV and 7.7 PeV.

128

135

136

137

138

61 **2** Data sample

The data used in this analysis were collected from May 2010 62 to May 2011 with 79 operational strings of IceCube. The 63 analysis was performed as a blind analysis, the selection 64 criteria to reject the background were developed using 10% 65 of the data ("burnsample"). This burnsample consists of 66 data uniformly distributed over the year to avoid biases 67 in muon background rate due to seasonal variations. The 68 burnsample livetime was 33 days. The numbers presented 69 here are based on the remaining 90% of the data, 317 days. 70 The main background for a search for cascade-like events 71 comes from cosmic ray muons with a faint track and a 72 single catastrophic energy loss from a bremsstrahlung. The 73 background of atmospheric muon events was simulated with 74 the air-shower program CORSIKA [8]. The main goal was 75 to simulate high energy muons that radiate bremsstrahlung 76 secondaries with energies that can mimic cascade events. 77 In this analysis, the CORSIKA background simulation 78 generated for the primary cosmic ray energy higher than 79 30 TeV per nucleon was used. A sample of 300 days of 80 atmospheric muon events in the energy range above 30 TeV 81 per nucleon was generated. 82

The signals in this analysis are v_e and v_{τ} from CC and 83 84 all neutrino flavors cascades from NC interactions. The all flavor neutrino events were simulated with the neutrino¹²¹ 85 generator ANIS [9] for energies from 1 TeV to 1 EeV at²² 86 the surface of the Earth with E^{-2} energy spectrum. Equal²³ 87 amounts of v and \bar{v} was produced. IceCube does not distin¹²⁴ 88 guish v from \bar{v} and in this paper v denotes the sum of v^{125} 89 and \bar{v} . In this analysis we used the flux normalization of 12690 127 signal events of 91

$$\Phi_{model} = 1.0 \times 10^{-8} (E/GeV)^{-2} GeV^{-1} s^{-1} sr^{-1} cm^{-2}.$$
(1)¹³⁰

The background from atmospheric neutrinos was esti¹³¹
 mated assuming the conventional [10] and prompt [11] flux¹³²
 contributions.

96 **3** Analysis

97 **3.1 Cascade reconstruction variables**

To isolate the cascade signal from muon background, dif_{t39} ferent selection criteria were applied. Among these were₁₄₀ simple quality criteria like the specific topology of cascade₁₄₁ like events, the development of the hit pattern in time, as₁₄₂ well as causal and likelihood criteria. 143

A widely utilized topology criterion for cascade analysis₁₄₄ 103 was provided by TensorOfInertia [2]. This reconstruc₁₄₅ 104 tion considered the hit-pattern as a rigid body, with the op₁₄₆ 105 tical modules as mass points with their charge equivalent₄₇ 106 to their mass. For this rigid body, the mass-eigenstates and 48 107 corresponding eigenvalues were calculated. The ratio of the149 108 highest eigenvalue and the sum of all three eigenvalues is150 109 a measure how spherical the hit-pattern is and thus can be₁₅₁ 110 used to separate cascade-like from track-like events. 111 152

To separate a cascade-like hit pattern, which is a station₁₅₃ ary source of light and a track, a moving source of light₁₅₄ the hits in the detector were projected along a track mov₁₅₅ ing through the detector with LineFitVelocity [12]. Fon₅₆ cascade-like events its value is much smaller then for track₁₅₇ like events and the identification of both hit patterns was₁₅₈ possible. 159

¹¹⁹ In the analysis chain the following likelihood reconstruc_{±60} ¹²⁰ tion algorithms were used: ACER [13], which is a determin_{±61}



Figure 1: Normalized TimeSplitPosition distributions for data (black points), sum of muon and atmospheric neutrino backgrounds (blue) and E^{-2} astrophysical v_e signal (magenta).

istic energy estimator, CascadeLlh [2], which uses probability density functions (pdfs) to perform a 4-dimensional fit, and Credo, which is more sophisticated algorithm that incorporates a model of light propagation in the ice, the full timing information and reconstructs the energy and direction of the incident neutrino.

The FillRatio was used to distinguish cascade-like events from muon-like tracks. Firstly, the mean distance between the vertex position and all hit DOMs in an event was calculated. Then, the ratio of number of hit DOMs to the total of all DOMs in the sphere of this mean radius was obtained. For a neutrino signal (cascade-like events) we expect this number to be close to one while for the track-like events this number would be uniformly distributed. This allows us to separate signal from background.

Another topology variable used in this analysis was TimeSplitPosition. Each event was split into two halves based on the charge-weighted mean time, and the cascade reconstruction was run on each half separately. Then, the difference TimeSplitPosition between reconstructed vertex positions for both halves was calculated. For the events consistent with a signal cascade hit pattern this number has a smaller value than for track-like events and allows separation of signal from background, as shown in Fig. 1. Figure 1 shows the normalized TimeSplitPosition distributions for data, Monte Carlo background and E^{-2} astrophysical v_e signal. The shape of the data distribution is nicely reproduced by the sum of muon and atmospheric backgrounds and represents the typical data-Monte Carlo shape agreement at different cut levels in the analysis presented here.

The ratio of maximum total charge on a single DOM in a given event and the total charge in this event MaxQTotRatio allowed the identification of the events, where most of the charge was recorded by a single DOM. These events might be created by a low energy muon having a catastrophic energy loss next to a DOM.

The variable DelayTime, defined as a minimum of the time difference between the first hit on a DOM and the time of the reconstructed vertex was also used. It allows the separation of a muon-track and cascade-like events as for the former this time difference is bigger than for the latter.







Figure 2: Schematic top view of IceCube with 79-strings. The green denotes the most outer layer of strings.

162 **3.2 Online filters**

To reduce the background coming from atmospheric muons and muon bundles several filters were applied to the data. The online filtering process begins at the South Pole with a trigger logic to suppress electronic noise and noise induced by radioactive processes of the detector itself.

The main physics trigger in IceCube is a "Simple Mul-168 tiplicity Trigger" (SMT) that requires photon signals in at 169 least 8 DOMs. The average trigger rate for the IceCube 79-170 string configuration was 1970 Hz. In the cascade online fil-171 ter the cuts on TensorOfInteria and LineFitVelocity 172 were applied to select cascade-like signal events from track-173 174 like background. The online filter reduced the data rate to 175 21 Hz, about a factor of 100 below the trigger rate. The cascade filter retained 75% of the v_e signal. After applying 176 the online filter, the data stream was transferred to the North 177 where more elaborate CascadeLlh and ACER cascade re-178 constructions were performed. 179

180 **3.3 Event selection**

204

Selection criteria to reject muon and atmospheric neutrinœ05
 backgrounds were developed. The Level3 filter retained₀₆
 events that fulfilled either a combined criterion of a cascade₂₀₇
 and track likelihood ratio LlhRatio as well as an energy₂₀₈
 dependent zenith angle cut or had a reconstructed ACER₂₀₉
 energy larger than 10 TeV.

Then the data stream was split into two branches: $fully^{210}$ 187 contained and partially contained events and each branch211 188 was analyzed separately. Only the fully contained events²¹² 189 selection criteria are described here but the partially con²¹³ 190 tained events were used to enhance the sensitivity of this214 191 analysis for neutrino events with energies E > 100 TeV. 215 192 The fully contained events were considered those with 216 193 both the reconstructed vertex and the first hit inside the₂₁₇ 194 most outer string layer of the detector, the green polygon₂₁₈ 195 in Fig. 2. In addition, we required that the first hit in the $_{219}$ 196 event occurred between ± 430 meters in depth and the 197 reconstructed Credo vertex position Z was between $\pm 450^{220}$ 198 meters in the detector. We rejected the event if the earlies t^{221} 199 hit occurred in the seven topmost DOMs. The FillRatio²²² 200 was calculated for this branch and only events with value²²³ 201 higher the 0.6 were retained. 202

At the Level4, further cuts were applied to reduce the225



Figure 3: Distribution of the reconstructed energy before final energy cut.



Figure 4: Model Rejection Factor (MRF) as a function of reconstructed energy.

background from atmospheric muons. Based on the time and position of the pulses in a given event, the events seen by 4 or more strings were selected. In the next step of Level4, we required that the reconstructed energy was higher that 10 TeV.

At Level5 we retained events with TimeSplitPosition smaller than 40 meters and rejected events with MaxQTotRatio bigger than 0.35. In addition, we required that the DelayTime was bigger than 100 ns.

Finally, using the Feldman-Cousins method [14], a cut on reconstructed energy (see Fig. 3) was optimized and used to suppress remaining muon and atmospheric neutrinos background. The Model Refection Factor (MRF) [15] was calculated as a function of reconstructed energy as shown in Fig. 4. The minimum of the MRF distribution was found at an energy of E=40 TeV and the energy cut was placed at this value. The energy resolution for an E^{-2} astrophysical spectrum for fully contained events is $\Delta(\log_{10}E_{\nu}) \sim 0.04$ and the vertex position resolution is ~ 4 meters.

The analysis aiming at partially contained astrophysical neutrino search has a poorer energy resolution of $\Delta(\log_{10}E_v) \sim 0.3$, and the vertex resolution of ~ 10 meters.





Figure 5: Effective area after the final event selection.

226 4 Results

The selection criteria rejected all of the CORSIKA events 227 and the conservative estimate on the number of cosmic-ray 228 muons at the final level was taken as an upper boundry at²⁷² 229 90% C.L. interval of 1.6 events. One burn sample data event²⁷³ 230 of 70 TeV reconstructed energy was retained. 274 231 From the analysis presented here 4.1 ± 0.2 (stat) v_e , 0.83 232 \pm 0.07 (stat) v_{μ} and 2.76 \pm 0.06 (stat) v_{τ} signal events for 233 an astrophysical flux defined in Eq. (1) are expected in 317^{75} 234 days (90% of the experimental data). Thereby, the predicted₂₇₆ 235

number of astrophysical v_{μ} events from CC interactions is₂₇₇ 0.31 ± 0.04 (stat), while from NC is 0.52 ± 0.05 (stat). 278

The expected number of atmospheric neutrino back₂₇₉ 238 ground events from v_e is 2.5 ± 0.2 (stat) +3.1-2.5 (syst) and₂₈₀ 239 from v_{μ} 1.8 ± 0.2 (stat) ± 0.6 (syst). The statistical uncer₂₈₁ 240 tainties come from the Monte Carlo statistics. The uncer282 241 tainties of the theoretical models in the predicted fluxes are₂₈₃ 242 dominating sources of systematic uncertainties for estimat₂₈₄ 243 ing atmospheric neutrino background. The uncertainty of_{85} 244 25% for conventional [10] and the factor of two for prompt₈₆ 245 flux [11] were assumed. These atmospheric background es₂₈₇ 246 timates include the neutrino events that would be accom₇₈₈ 247 panied by a muon bundle [17] and therefore removed by₂₈₉ 248 the analysis selection cuts. The estimated background could₉₀ 249 hence be lowered by a factor of ~ 2 . 250 291

Figure 5 shows the effective area versus neutrino energy₂₉₂ after all cuts applied. The Glashow resonance [16] contribu₂₉₃ tion is clearly visible for v_e . The effective areas for v_e and₂₉₄ v_{τ} are higher than for v_{μ} as this analysis was optimized fo₅₉₅ cascades and removed muon tracks. 296

The comparison of the all neutrino flavor effective area₂₉₇ for combined fully and partially contained analyses with 79₂₉₈ IceCube strings and the cascade search with the 59-string₂₉₉ IceCube configuration is shown in Fig. 6. The effective area₈₀₀ for the 79-string configuration is bigger than for a smaller₅₀₁ detector, as expected. 302

The sensitivity for the diffuse all flavor flux of extrater₃₀₃ 262 restrial neutrino signal, defined as the average flux upper₃₀₄ 263 limit at 90% C.L. in the absence of signal was calculated₃₀₅ 264 and resulted in 2.3×10^{-8} GeV s⁻¹ sr⁻¹ cm⁻² for the all-265 flavor neutrino energies between 42 TeV and 6 PeV. No 266 systematic uncertainties were taken into account. Including 267 partially contained events increases the sensitivity to 1.8 268 $\times 10^{-8}$ GeV s⁻¹ sr⁻¹ cm⁻² for all-flavor neutrino events 269 with energies between 44 TeV and 7.7 PeV. The obtained 270 result is more stringent than the expected upper limits from 271



Figure 6: Comparison of effective area for sum of all flavor neutrinos (open squares) for the analysis presented here and the cascade neutrino search with IC59 string configuration [5] (filled circles).

previous IceCube cascade analyses with smaller sized detector configurations [2, 4, 5]. The systematic uncertainties are currently being evaluated.

References

- [1] A. Achterberg *et al.*, Astropart. Phys. **26** (2006) 155 doi: 10.1016/j.astropartphys.2006.06.007.
- [2] R. Abbasi *et al.*, Phys. Rev. **D84** (2011) 072001 doi:10.1103/PhysRevD.84.072001.
- [3] E. Middell *et al.*, Proceedings of the 32nd ICRC, 2011 Included in arXiv:astro-ph/1111.2736.
- [4] S. Hickford, S. Panknin *et al.*, Proceedings of the 32nd ICRC, 2011 Included in arXiv:astro-ph/1111.2736.
- [5] IceCube Collaboration, paper 0662 these proceedings.
- [6] M.G. Aartsen et al., arXiv:astro-ph/1304.5356.
- [7] IceCube Collaboration, paper 0650 these proceedings.
- [8] D. Heck et al., Tech. Rep. FZKA (1998) 6019.
- [9] A. Gazizov, M. Kowalski, Computer Physcis Communications, Vol 172 (2005) 203; arXiv:astro-ph/0406439.
- [10] M. Honda *et al.* Phys. Rev. **D75** (2007) 043006 doi: 10.1103/PhysRevD.75.043006.
- [11] R. Enberg, M.H. Reno, I. Sarcevic, Phys. Rev. D78 (2008) 043005 doi: 10.1103/PhysRevD.78.043005.
- [12] J. Ahrens *et al.*, Phys. Rev. **D67** (2003) 012003 doi: 10.1103/PhysRevD.67.012003.
- [13] M. D'Agostino, Ph.D. thesis, University of California, Berkeley (2009), arXiv:astro-ph/0910.2555.
- [14] G.J. Feldman, R.D. Cousins, Phys. Rev. D57 (1998) 3878 doi: 10.1103/PhysRevD.57.3873.
- [15] G.C. Hill, K. Rawlins, arXiv:astro-ph/0209350.
- [16] S. L. Glashow, Phys. Rev. 118 (1960) 316 doi: 10.1103/PhysRev.118.316.
- [17] S. Schönert *et al.*, Phys. Rev. **D79** (2009) 043009 doi: 10.1103/PhysRevD.79.043009.