# dark matter searches with IceCube

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- · IceCube is complete ! (last string down on Dec 18th 2010)
  - taking data with 86 strings since May 2011
- · Results from 2008/09 (22/40 strings) are being released
- · IceCube reaches the Galactic Center and Halo

opens possibility for a multi-wavelength approach to dark matter searches

- · Analysis on 79 strings in the pipeline
- New statistical techniques being developed: global model fits using all available data
- IceCube searches for DM competitive with direct searches. Natural extension to collider searches
- The low-energy extension, DeepCore, deployed and taking data

### the IceCube neutrino telescope

7 last strings installed during Dec 2010

86 strings and 80 IceTop tanks

### complete!



IceTop: Air shower detector

80 stations/2 tanks each threshold ~ 300 TeV

1450 m

Inice array: 80 Strings 60 Optical Modules 17 m between Modules 125 m between Strings v threshold ≤100 GeV 2450 m

> DeepCore array: 6 additional strings 60 Optical Modules 7/10 m between Modules 72 m between Strings v threshold ~10 GeV

### the IceCube Digital Optical Module

### Each DOM is an autonomous data collection unit



- PMT: Hamamatsu, 10"
- Digitizers:
  - ATWD: 3 channels. Sampling 300MHz, capture 400 ns

<u>FADC</u>: sampling 40 MHz, capture 6.4  $\mu$ s

Dynamic range 500pe/15 nsec, 25000 pe/6.4  $\mu$ s



- Flasher board: 12 controllable LEDs at 00 or 450

Clock stability:  $10-10 \approx 0.1$  nsec / sec Synchronized to GPS time every  $\approx 5$  sec at 2 ns precision

### dark matter candidates



#### <u>WIMPS</u>

- Arise in extensions of the Standard Model
- Assumed to be stable: relics from the Big Bang
- weak-type Xsection gives needed relic density

$$\Omega_{\delta} h^2 \approx \frac{10^{-27}}{\langle \sigma_{ann} v \rangle_{fr}} \,\mathrm{cm}^3 \,\mathrm{s}^{-1}$$

- mass from few GeV to few TeV
- **MSSM** candidate: lightest neutralino,

 $\chi \square_{1}^{0} = N_{1}B + N_{2}W^{3} + N_{3}H^{0}_{1} + N_{4}H^{0}_{2}$ 

- UED: lightest 'rung' in the Kaluza-Klein ladder

#### **SIMPZILLAS**

- Non-thermal, non-weakly interacting stable relics

- A wealth of candidates from different theoretical models:
  - dark baryons (primordial nucleosynthesis constraints)
  - MACHOs BHs, neutron stars, white/brown dwarfs... (microlensing constraints)
  - neutrinos (mass constraint)
  - primordial Black Holes (cosmological constraints)
  - Weakly Interacting Massive Particles (LSPs from "x"MSSM, Kaluza-Klein modes...)
  - Non-weakly Interacting Supermassive particles (Simpzillas)
  - axions (too light+astrophysical constrains)
  - many others
- ... + (alternative gravity theories)

# DM-induced SM particles:

 $\kappa\kappa, \chi\chi, SS \rightarrow \left\{ \begin{array}{c} q\overline{q} \\ \ell^{+}\ell^{-} \\ W, Z, H \end{array} \right\} \rightarrow \nu, \gamma, e^{+}e^{-}, \overline{p}$ 

Kaluza-Klein modes an additional useful channel:

 $\kappa\kappa \to \, \nu\nu$ 

signature:

v excess over background from Sun/Earth/Galactic Halo/nearby galaxies

A lot of physics uncertainties involved:

- relic density calculations
- DM distribution in the halo
- velocity distribution
- $\chi$ ,K,S properties (MSSM/UED...)
- interaction of  $\chi$ ,K,S with matter (capture)
- self interaction (annihilation)

### indirect searches for dark matter

Look at objects where dark matter can have accumulated gravitationally over the evolution of the Universe

# Sun, Earth, Galactic Halo/Center, dwarf spheroids



Atmospheric muons ~O(10<sup>10</sup>) events/year (downwards)

Atmospheric neutrinos ~O(10<sup>4</sup>) events/year (all directions)

Triggered data still dominated by atmospheric muons

Reject misreconstructed atmospheric muon background through event and track quality parameters

Use of **linear cuts** and/or multivariate methods to extract irreducible atmospheric neutrino background (Neural Nets, Support Vector Machines, Boosted Decision Trees)

DM searches directional: good additional handle on event selection

### $\rightarrow$ distribution-shape analysis

(allow for a higher background contamination)

sequential cuts



Solid angle to the Sun  $\psi$  ( $\rho\alpha\delta$ )

## analysis strategies in neutrino telescopes

$$\begin{array}{c} \mathbf{N}_{\text{data'}}, \mathbf{N}_{\text{bck}} \\ \mathbf{\Psi}_{\text{data'}}, \mathbf{\Psi}_{\text{bck}} \\ \hline \rightarrow \mathbf{N}_{90} \\ \hline \mathbf{F}_{\nu\mu} \leq \frac{\mathbf{N}_{90}}{V_{\text{eff}} \cdot t} \\ \hline \mathbf{Experimentally obtained quantity:} \\ \text{allowed number of signal events still compatible with background, at 90% confidence level} \\ \hline \mathbf{F}_{\nu\mu}(m_{\chi}) = \mathbf{F}_{A} \cdot \frac{1}{4\pi R_{\oplus}^{2}} \int_{0}^{m_{\chi}} \sum B_{\chi\chi \rightarrow \chi} \left( \frac{dN_{\nu}}{dE_{\nu}} \right) \\ \times \sigma_{\nu+N \rightarrow \mu} + \dots \left( E_{\nu} | E_{\mu} \geq E_{\text{thr}} \right) \rho_{N} dE_{\nu} \\ \hline \end{array}$$

$$\begin{array}{c} \text{Use model to convert} \\ \text{to a muon flux} \\ \hline \phi & \phi_{\mu}(E_{\mu} \geq E_{\text{thr}}) = \frac{\Gamma_{A}}{4\pi D_{\odot}^{2}} \int_{E_{\text{thr}}}^{\infty} dE_{\mu} \frac{dN_{\mu}}{dE_{\mu}} \\ \hline \\ \psi & \phi_{\mu}(E_{\mu} \geq E_{\text{thr}}) = \frac{\Gamma_{A}}{4\pi D_{\odot}^{2}} \int_{E_{\text{thr}}}^{\infty} dE_{\mu} \frac{dN_{\mu}}{dE_{\mu}} \\ \hline \\ \psi & \phi_{\mu}(E_{\mu} \geq E_{\text{thr}}) = \frac{\Gamma_{A}}{4\pi D_{\odot}^{2}} \int_{E_{\text{thr}}}^{\infty} dE_{\mu} \frac{dN_{\mu}}{dE_{\mu}} \\ \hline \\ \psi & \phi_{\mu}(E_{\mu} \geq E_{\text{thr}}) = \frac{\Gamma_{A}}{4\pi D_{\odot}^{2}} \int_{E_{\text{thr}}}^{\infty} dE_{\mu} \frac{dN_{\mu}}{dE_{\mu}} \\ \hline \\ \psi & \phi_{\mu}(E_{\mu} \geq E_{\text{thr}}) = \frac{\Gamma_{A}}{4\pi D_{\odot}^{2}} \int_{E_{\text{thr}}}^{\infty} dE_{\mu} \frac{dN_{\mu}}{dE_{\mu}} \\ \hline \\ \psi & \phi_{\mu}(E_{\mu} \geq E_{\text{thr}}) = \frac{\Gamma_{A}}{4\pi D_{\odot}^{2}} \int_{E_{\text{thr}}}^{\infty} dE_{\mu} \frac{dN_{\mu}}{dE_{\mu}} \\ \hline \\ \psi & \phi_{\mu}(E_{\mu} \geq E_{\text{thr}}) = \frac{\Gamma_{A}}{4\pi D_{\odot}^{2}} \int_{E_{\text{thr}}}^{\infty} dE_{\mu} \frac{dN_{\mu}}{dE_{\mu}} \\ \hline \\ \psi & \phi_{\mu}(E_{\mu} \geq E_{\text{thr}}) = \frac{\Gamma_{A}}{4\pi D_{\odot}^{2}} \int_{E_{\text{thr}}}^{\infty} dE_{\mu} \frac{dN_{\mu}}{dE_{\mu}} \\ \hline \\ \psi & \phi_{\mu}(E_{\mu} \geq E_{\text{thr}}) = \frac{\Gamma_{A}}{4\pi D_{\odot}^{2}} \int_{E_{\text{thr}}}^{\infty} dE_{\mu} \frac{dN_{\mu}}{dE_{\mu}} \\ \hline \\ \psi & \phi_{\mu}(E_{\mu} \geq E_{\text{thr}}) = \frac{\Gamma_{A}}{4\pi D_{\odot}^{2}} \int_{E_{\text{thr}}}^{\infty} dE_{\mu} \frac{dN_{\mu}}{dE_{\mu}} \\ \hline \\ \psi & \phi_{\mu}(E_{\mu} \geq E_{\text{thr}}) = \frac{\Gamma_{A}}{4\pi D_{\odot}^{2}} \int_{E_{\text{thr}}}^{\infty} dE_{\mu} \frac{dN_{\mu}}{dE_{\mu}} \\ \hline \\ \psi & \phi_{\mu}(E_{\mu} \geq E_{\text{thr}}) = \frac{\Gamma_{A}}{4\pi D_{\odot}^{2}} \int_{E_{\text{thr}}}^{\infty} dE_{\mu} \frac{dN_{\mu}}{dE_{\mu}} \\ \hline \\ \psi & \phi_{\mu}(E_{\mu} \geq E_{\text{thr}}) = \frac{\Gamma_{A}}{4\pi D_{\odot}^{2}} \int_{E_{\text{thr}}}^{\infty} dE_{\mu} \frac{dN_{\mu}}{dE_{\mu}} \\ \hline \\ \psi & \phi_{\mu}(E_{\mu} \geq E_{\mu}) = \frac{\Gamma_{A}}{4\pi D_{\oplus}^{2}} \int_{E_{\text{thr}}}^{\infty} dE_{\mu} \frac{dN_{\mu}}{dE_{\mu}} \\ \hline \\ \psi & \phi_{\mu}(E_{\mu} \geq E_{\mu}) = \frac{\Gamma_{A}}{4\pi D_{\oplus}^{2}} \int_{E_{\mu}(E_{\mu} \otimes E_{\mu}) } \\ \psi$$

### searches from the Sun: results

#### IceCube results from 1065 days of livetime between 2001-2008

#### 90% CL muon flux limit from the Sun

(compared to MSSM scans)

#### 90% CL neutralino-p Xsection limit

(compared to MSSM scans)



(particle physics and solar model)

Assume (ie. model dependent) effective quark-DM interaction,

 $\lambda^2/\Lambda^2 (\overline{q}\gamma_5\gamma_\mu q)(\overline{\chi}\gamma_5\gamma^\mu \chi)$ 

and look for monojets in  $p\overline{p}$  collisions,

 $p\overline{p} \rightarrow \chi \overline{\chi} + jet$ 

(as opposed to the SM process  $pp \rightarrow Z+jet$  and  $pp \rightarrow W+jet$ )

Constrains from monojet searches at the TeVatron:



#### 90% CL neutralino-p Xsection limit



Universal Extra Dimensions:



 $n=1 \rightarrow Lightest Kaluza-Klein mode, B^1$ 

good DM candidate

90% CL LKP-p Xsection limit vs LKP mass



### searches from the Sun: Superheavy DM

<u>SIMPZILLAS (</u>Superheavy DM)

- Produced **non-thermally** at the end of inflation through vacuum quantum fluctuations or decay of the inflaton field

- strong Xsection (simply means non-weak in this context)

- m from  $\sim 10^4$  GeV to  $10^{18}$  GeV (no unitarity limit since production non thermal)

 $v, \overline{q}$ 



$$N_s(m_{\rm x},\sigma_{\rm XN}) = N_t \cdot BR_{\rm w} \cdot \Gamma_A(m_{\rm x},\sigma_{\rm XN}) \cdot T \cdot \int \frac{dN_\nu}{dE} A_{eff} dE$$

90% CL simpzilla-p Xsection limit vs simpzilla mass



Extend the search to the southern hemisphere by selecting starting events

 $\rightarrow$  Veto background through location of interaction vertex

- muon background: downgoing, no starting track
- WIMP signal: interaction vertex within detector volume

#### work in progress







#### **Starting Track Interaction Vertex**

### searches from the Sun: secluded dark matter

#### Secluded dark matter

 $\mathcal{L} = \mathcal{L}_{SM} + \mathcal{L}_{WIMP} + \mathcal{L}_{mediator}$ 

DM annihilates into mediator,  $\chi\chi \rightarrow \phi\phi \rightarrow SM$ with  $m_{\phi} = O(GeV)$ 

 $\phi$  is long lived, escapes the Sun and decays into  $\mu^+\mu^-$  in or near the detector

 $\rightarrow$  signature: two closely separated muon tracks (~ 1m)

look for stopping pairs of tracks in order to further reduce the background.

work in progress





### searches from the Sun: neutrino energies at the detector



: Indirect dark matter searches from the **Sun** are a low-energy analysis in neutrino telescopes: even for the highest DM masses, we do not get muons above few 100 GeV

Not such effect for the Earth and Halo (no v energy losses in dense medium)

Explaining Fermi/PAMELA data in terms of dark matter favors boost in DM annihilation Xsection

If the dark matter annihilation rate is enhanced, the timescale for equilibrium diminishes  $\rightarrow$  flux of annihilation products can be much larger than away from equilibrium.

→ an enhanced annihilation Xsection can produce a detectable neutrino flux from the center of the Earth (while not enhancing the Solar flux) (C. Delaunay, P. J. Fox and G. Perez, JHEP 0905 , 099 (2009)).

Using the atmospheric neutrino measurement of IceCube-40, model-independent limits on boost factors can be set







### searches from the Galactic Halo





### searches from the Galactic Center



Look for an excess of events in the on-source region w.r.t. the off-source

on-source region below the horizon: need to veto downgoing  $\mu$ s.

Use central strings of detector as fiducial volume, surrounding layers as veto. Only from IC40 this is possible.

IC40: observed on-source: 798842 evts observed off-source: 798819 evts

Same strategy as in the galactic halo analysis:



### IceCube22/40 results from DM searches from the Galaxy



 line thickness in IC22 results reflects uncertainty due to halo profile

 no energy loss of secondaries before decay: harder spectra than in Sun for same annihilation channel



IC22: Phys. Rev. 84, 022004 (2011) IC40: in preparation Galactic Center results in the framework of satellite searches

· multi-wavelength approach to dark matter searches:

IceCube results in the context of Pamela and Fermi anomaly



### searches from nearby dwarf galaxies: strategy

- dwarf galaxies: high mass/light ratio
- $\rightarrow$  high concentration of DM in the halos
- known location. Distributed both in the north and southern sky.
  - Point-like search techniques: stacking
  - known distance -> determination of absolute annihilation rate if a signal is detected
- same expected neutrino spectra as for the galactic center/halo

IceCube analysis in progress

Same strategy as in the galactic halo analysis:

$$\frac{d\Phi_j(\Delta\Omega, E_j)}{dE_j} = \frac{\langle \sigma v \rangle}{2m_\chi^2} \frac{dN_j}{dE_j} J(\Delta\Omega)$$



10<sup>4</sup> WIMP mass (GeV)

### the DeepCore infill

• Aim: lower energy threshold through a denser core in the center of the IceCube array

 6 additional strings of 60 high quantum efficiency PMTs

denser instrumentation:
7 m DOM vertical spacing (17m in IceCube),

72 m inter string spacing (125m in IceCube)



full sky sensitivity using IceCube surrounding strings as a veto:

375m thick detector veto: three complete IceCube string layers surround DeepCore

—> access to southern hemisphere, galactic center and all-year Sun visibility

Preliminary studies show 10<sup>3</sup> background rejection with 99% signal efficiency possible at filter level



### DeepCore veto capabilities



- $\cdot$  we have 1 km<sup>3</sup> of ice instrumented with optical modules
- · we can detect flavours (muon tracks,  $e/\tau$  cascades)
- · we can define through-going, starting and contained tracks
- · we cover a wide neutrino energy range, from few tens GeV to PeV
- we can look at all the sky (at once and continuously)

..... if you have a model of exotic physics that involves neutrinos, we can probe it